Abstract

In this thesis the global Standard Model (SM) fit to the electroweak precision observables is revisited with respect to newest experimental results. Various consistency checks are performed showing no significant deviation from the SM. The Higgs boson mass is estimated by the electroweak fit to be $M_H = 94^{+30}_{-24}$ GeV without any information from direct Higgs searches at LEP, Tevatron, and the LHC and the result is $M_H = 125^{+8}_{-10}$ GeV when including the direct Higgs mass constraints. The strong coupling constant is extracted at fourth perturbative order as $\alpha_s(M_Z^2) = 0.1194 \pm 0.0028$ (exp) ± 0.0001 (theo). From the fit including the direct Higgs constraints the effective weak mixing angle is determined indirectly to be $\sin^2 \theta_{\rm eff}^\ell = 0.23147^{+0.0012}_{-0.0010}$. For the W mass the value of $M_W = 80.360^{+0.012}_{-0.011}$ GeV is obtained indirectly from the fit including the direct Higgs constraints.

The electroweak precision data is also exploited to constrain new physics models by using the concept of oblique parameters. In this thesis the following models are investigated: models with a sequential fourth fermion generation, the inert-Higgs doublet model, the littlest Higgs model with T-parity conservation, and models with large extra dimensions. In contrast to the SM, in these models heavy Higgs bosons are in agreement with the electroweak precision data.

The forward-backward asymmetry as a function of the invariant mass is measured for $pp \rightarrow Z/\gamma^* \rightarrow e^+e^-$ events collected with the ATLAS detector at the LHC. The data taken in 2010 at a center-of-mass energy of $\sqrt{s} = 7$ TeV corresponding to an integrated luminosity of 37.4 pb⁻¹ is analyzed. The measured forward-backward asymmetry is in agreement with the SM expectation. From the measured forward-backward asymmetry the effective weak mixing angle is extracted as $\sin^2 \theta_{\rm eff}^{\ell} = 0.2204 \pm 0.0071 \, ({\rm stat}) \, {}^{+0.0039}_{-0.0044} \, ({\rm syst})$. The impact of unparticles and large extra dimensions on the forward-backward asymmetry at large momentum transfers is studied at generator level.

Tests of the Electroweak Standard Model and Measurement of the Weak Mixing Angle with the ATLAS Detector

> Dissertation zur Erlangung des Doktorgrades des Department Physik der Universität Hamburg

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CONTENTS

5	The	e Experimental Set-Up 4			
	5.1	The Large Hadron Collider	47		
	5.2	The ATLAS Detector			
		5.2.1 The Coordinate System	49		
		5.2.2 The Magnet System	50		
		5.2.3 The Inner Detector	51		
		5.2.4 The Calorimetry	52		
		5.2.5 The Muon Spectrometer	56		
	5.3	Electron Reconstruction and Identification	57		
		5.3.1 Electron Reconstruction	57		
		5.3.2 Electron Identification	58		
	5.4	The ATLAS Trigger System	60		
		5.4.1 The Level-1 Trigger System	62		
		5.4.2 The High Level Trigger System	64		
		5.4.3 Electron Trigger Selection	64		
		5.4.4 The Trigger Menu and Configuration	66		
	5.5	Data Quality and Luminosity Determination	67		
		5.5.1 Electron Object Quality	69		
	5.6	Monte Carlo Generation and Detector Simulation	70		
c	171.	the minute Defense of	79		
0	Elec	L1 Derformance	73		
	0.1	6.1.1. Li Tuimme Dates	79		
		6.1.1 L1 Ingger Rates	75 75		
		0.1.2 L1 Ingger Emclency	() 77		
	0.0		11		
	6.2	HLT Performance	80		
7	Di-l	Electron Event Selection	85		
	7.1	Selection of $Z/\gamma^* \to ee$ candidates	85		
		7.1.1 Corrections and Efficiencies	87		
		7.1.2 Yield of $Z/\gamma^* \to ee$ candidates	89		
	7.2	Background Contamination	90		
		7.2.1 Data-Driven QCD Background Estimation	91		
	7.3	Di-Electron Distributions	95		
8	Mo	summer of $A_{}$ and $\sin^2 \theta^{\ell}$	101		
0	8.1	Analysis Strategy	101		
	8.2	Charge Mis_Identification			
	8.2	Consign Mis-infinite Margarety at Detector Lovel			
	0.0 8.4	Determination of the Effective Week Mixing Angle	107		
	85	Unfolded Forward Backward Asymmetry	112		
	0.0	0 monded Forward-Dackward Asymmetry	110 119		
		0.0.1 mass migration and Unarge mis-identification	119		

Contents

1 Introduction					
2	Theoretical Background				
	2.1	2.1 The Standard Model of Particle Physics			
		2.1.1	Problems of the Standard Model	6	
		2.1.2	The Effective Weak Mixing Angle	6	
	2.2	The D	Orell–Yan Process at the LHC	8	
3 The Global Fit of the Electroweak Standard Model				13	
	3.1 Statistical Aspects				
	3.2	3.2 Standard Model Predictions			
	3.3	Exper	imental Inputs	16	
	3.4 Fit Results		esults	18	
		3.4.1	Higgs Mass Constraints	21	
		3.4.2	Determination of the Strong Coupling	23	
		3.4.3	Determination of the Top Mass	23	
		3.4.4	Determination of the W Mass	25	
		3.4.5	Determination of the Effective Weak Mixing Angle	25	
		3.4.6	Two-Dimensional Scan of the W and Top Mass	26	
4	Cor	nstrain	ts on New Physics from the Electroweak Fit	29	
	4.1	Conce	pt of Oblique Parameters	29	
	4.2	.2 Experimental Constraints on Oblique Parameters .3 Constraints on New Physics Models		31	
	4.3			31	
		4.3.1	Inert-Higgs Doublet Model	34	
		4.3.2	Models with a Sequential Fourth Fermion Generation $\hfill \ldots \ldots \ldots$.	36	
		4.3.3	Littlest Higgs with T-Parity Conservation	39	
		4.3.4	Models with Large Extra Dimensions $\hfill \ldots \ldots \ldots \ldots \ldots \ldots$	43	

CONTENTS

CHAPTER 3	3
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The Global Fit of the Electroweak Standard Model

Precision measurements, in line with accurate theoretical predictions, allow us to probe physics at much higher energy scales than the masses of the particles directly involved in experimental reactions by exploiting contributions from quantum loops. In case of the SM, unknown parameters of the SM (*e.g.* the Higgs boson mass) can be determined from multi-parameter fits. The analysis presented in this chapter relies on the Gfitter framework [1] and the presentation follows closely (though not identically) the presentation in the publications [1, 34]. This analysis has benefited from the enormous work in the past which has been done for the calculation of the electroweak precision observables. During this effort various software packages have been developed predicting the electroweak precision observables within the SM: ZFITTER [35, 36], TOPAZO [37, 38], LEPTOP [39, 40], and GAPP [41, 42] (see also the review [43]). Electroweak SM fits are also routinely performed by the LEP Electroweak Working Group [44] and for the electroweak review of the Particle Data Group [42].

This chapter is organized as follows. First, the statistical aspects of the fits as implemented in the Gfitter framework and the theoretical predictions of the electroweak observables are discussed. The experimental data used in the analysis is introduced. Especially, the treatment of the information from direct Higgs searches is explained in detail. The chapter concludes with the presentation of various fit results. Among them, constraints on the Higgs mass, a determination of the strong coupling, and indirect determinations of M_W , m_t , and $\sin^2 \theta_{\text{eff}}^\ell$ are shown.

3.1 Statistical Aspects

The statistical analysis is performed with the Gfitter framework, which adopts a least-square like notation. The test statistics is defined as

$$\chi^2(y_{\rm mod}) \equiv -2\ln\mathcal{L}(y_{\rm mod})\,,\tag{3.1}$$

where the likelihood function (\mathcal{L}) depends on the free parameter (y_{mod}) of the physics model. The likelihood function of a parameter with its central measured value (x_0) , positive (negative)

		8.5.2	Correction for Acceptance and incorrect Quark Direction	114							
		8.5.3	Unfolding Procedure	117							
		8.5.4	Results	119							
	8.6	Impac	t of Models Beyond the SM	119							
		8.6.1	Large Extra Dimension	119							
		8.6.2	Unparticles	121							
		8.6.3	Conclusion	123							
9	Conclusions										
A	Number of Forward and Backward Events										
в	$_{\rm S}$ Systematic Uncertainties of Unfolded ${ m A_{FB}}$										
Bi	Bibliography										

3.2 Standard Model Predictions

3.2 Standard Model Predictions

The SM predictions for the electroweak observables measured by the LEP, SLC, and Tevatron experiments are implemented as a function of the floating fit parameters M_Z , M_H , m_t , $\overline{m_b}$, $\overline{m_c}$, $\Delta \alpha_{had}^5(M_Z^2)$, and $\alpha_s(M_Z^2)$. The predictions of the effective weak mixing angle and the W mass are the most important ones in order to constrain the mass of the Higgs boson. For the W mass eq. (6) and the coefficients of eq. (8) from [50] have been implemented. The calculation contains the full two-loop and leading beyond-two-loop corrections. The implementation of the effective weak mixing angle follows the full two-loop and leading beyond-two-loop computation of [51–53]. The asymmetry parameters can be computed by the effective weak mixing angle via

$$A_f = 2 \frac{g_V^f / g_A^f}{1 + (g_V^f / g_A^f)^2}, \qquad (3.5)$$

where

$$\frac{g_V^f}{g_A^f} = 1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f.$$
(3.6)

The forward-backward asymmetry where the superscript '0' indicates that the observed values have been corrected for radiative effects and photon exchange, can be determined from the asymmetry parameters as follows

$$A_{\rm FB}^{0,f} = \frac{3}{4} A_e A_f \,. \tag{3.7}$$

Unlike the asymmetry parameters, the partial widths of the Z boson are defined inclusively, *i.e.* they contain all real and virtual corrections. They can be computed by¹

$$\Gamma_Z^f = \frac{G_F M_Z^3}{6\sqrt{2\pi}} \left[(g_A^f)^2 \left((1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f) R_V^f + R_A^f \right) + \delta_{Im(\kappa^f)}^f R_V^f \right] + \Delta_{\text{EW/QCD}}^f, \qquad (3.8)$$

where G_F is the Fermi constant. For the analysis presented in this thesis, the vector and axialvector couplings $(g_A^f \text{ and } g_V^f)$ are implemented using the parametrizations of [55–58], which are computed at one-loop level and partly at two-loop level for $\mathcal{O}(\alpha \alpha_s)$.² To account for a different Higgs mass dependence of the parametrizations and ZFITTER at large Higgs masses $(M_H \gtrsim 500 \text{ GeV})$, a quadratic correction is added to $(\Delta T_Z)_{\rm SM}$ of eq. (23b) in [58] (for further explanation of the *T* parameter see chapter 4)

$$(\Delta T_Z)_{\rm SM,Corr} = \begin{cases} 0 & \text{if } M_H \le 200 \text{ GeV}, \\ 0.00764(x_h - x_{h,200}) - 0.112(x_h - x_{h,200})^2 & \text{if } M_H > 200 \text{ GeV}, \end{cases}$$
(3.9)

where $x_h = \log(M_H/100 \text{ GeV})$ and $x_{h,200} = \log(200/100).^3$ This correction factor does not affect the electroweak SM fit, but it slightly influences the constraints on new physics parameters (*cf.* chapter 4). The term $\delta_{I_m(\kappa f)}^f$ in eq. (3.8) represents the corrections from the imaginary part

3. THE GLOBAL FIT OF THE ELECTROWEAK STANDARD MODEL

Gaussian error σ_{Gauss}^+ (σ_{Gauss}^-), and positive (negative) theoretical error σ_{theo}^+ (σ_{theo}^-), for a given set of y_{mod} parameters and the theoretical prediction $f(y_{\text{mod}})$ is given by¹

$$-2\ln\mathcal{L}(y_{\rm mod}) = \begin{cases} 0, & \text{if: } -\sigma_{\rm theo}^{-} \leq f(y_{\rm mod}) - x_0 \leq \sigma_{\rm theo}^{+}, \\ \left(\frac{f(y_{\rm mod}) - (x_0 + \sigma_{\rm theo}^{+})}{\sigma_{\rm Gauss}^{-}}\right)^2, & \text{if: } f(y_{\rm mod}) - x_0 > \sigma_{\rm theo}^{+}, \\ \left(\frac{f(y_{\rm mod}) - (x_0 - \sigma_{\rm theo}^{-})}{\sigma_{\rm Gauss}^{-}}\right)^2, & \text{if: } x_0 - f(y_{\rm mod}) > \sigma_{\rm theo}^{-}. \end{cases}$$
(3.2)

Theoretical uncertainties are treated according to the *R*Fit scheme [45, 46], *i.e.* the theoretical prediction can freely vary within the range of the theoretical uncertainty without contributing to the χ^2 estimator. The final test statistics of the global fit is defined as the sum over all $-2 \ln \mathcal{L}(y_{\rm mod})$ contributions from each observable. Correlations between measurements are considered properly in the likelihood function.

In addition, it is possible to introduce dependencies among parameters in Gfitter, which can be used to parametrize correlations due to common systematic errors, or to rescale parameter values and errors with newly available results for parameters on which other parameters depend (rescaling mechanism).

For the parameter estimation the offset-corrected test statistics is used

$$\Delta \chi^2(y_{\text{mod}}) = \chi^2(y_{\text{mod}}) - \chi^2_{\text{min}}(y_{\text{mod}}), \qquad (3.3)$$

where $\chi^2_{\min}(y_{\text{mod}})$ is the absolute minimum of the test statistics. The minimum value of $\Delta \chi^2$ is zero, by construction. This ensures that, consistent with the assumption that the model is correct, exclusion confidence levels (CL) equal to zero are obtained when exploring the y_{mod} space.² The CL is computed for a Gaussian problem by

$$CL = 1 - Prob(\Delta \chi^2, n_{dof}), \qquad (3.4)$$

where $n_{\rm dof}$ is the number of degrees of freedom of the offset-corrected $\Delta \chi^2$. In case of a non-Gaussian problem a toy Monte Carlo analysis is required to estimate the CL. For the electroweak SM fit no significant deviations between the toy Monte Carlo analysis and eq. (3.4) are observed [1].

The p-value is an estimator for the goodness of the fit. It quantifies the probability of wrongly rejecting the theoretical hypothesis. For a Gaussian problem the p-value is compute by $\text{Prob}(\Delta\chi^2, n_{\text{dof}})$. In case of the electroweak fit, this naive p-value determinations have been confirmed with Monte Carlo toy experiments [1, 47].

In Gfitter the minimization of the test statistics is performed by TMinuit [48]. In addition, more involved global minima finders are used: Genetic Algorithm and Simulated Annealing, which are available with the TMVA [49] package in ROOT [2].

¹See for instance [54] and [55–58].

²The above mentioned Gfitter publications use the implementations from ZFITTER [35, 36], which contains up to two-loop electroweak corrections [35, 36, 43, 54, 59–66] and all known QCD corrections [35, 36, 67].

³The coefficients are determined by comparing the values of the T parameter computed with eq. (23b) of [58] and with the Fortran ZFITTER package [35, 36] (version 6.42 [68]).

¹The central value x_0 corresponds to the value with the largest likelihood, which is not necessarily equal to the arithmetic average in case of asymmetric errors.

²Throughout this thesis the term confidence level denotes 1 minus the p-value of a given $\Delta \chi^2$ (or χ^2) test statistics, and is hence a measure of the exclusion probability of a hypothesis. This is not to be confounded with a confidence interval, which expresses an inclusion probability.

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141

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